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S**EFFECT OF VARIABLE THERMAL PROPERTIES  
ON GUN TUBE HEATING****NATHAN GERBER  
MARK BUNDY****JULY 1992**

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The authors would like to thank Mr. Paul Conroy of the Interior Ballistics Division, Ballistic Research Laboratory (BRL), for supplying us with an interior ballistics prediction of the in-bore gas temperature and heat transfer coefficient, as a function of position and time, for an M256 120-mm gun barrel firing DM13 ammunition. These data were used as input in our computations. We also wish to thank Mr. Joe Cox, Benet Weapons Laboratory, for providing us with the measured values of specific heat as a function of temperature for M256 gun barrel steel. These temperature-dependent data were a critical part of our nonlinear analysis. And finally, we are indebted to Mr. Jim Bradley, Launch and Flight Division, BRL, for his adept skills at programming our mathematical formulation into an efficient FORTRAN code.

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## 1. INTRODUCTION

Gun tube heating from multiple firings continues to be a subject of concern to ordnance engineers. A number of investigators have modeled the heating. Some recent publications include: Artus and Hasenbein (1989); Chandra and Fisher (1989a, 1989b); Talley (1989a, 1989b); Rapp (1990); Chandra (1990); Gerber and Bundy (1991); and Conroy (1991). In these studies, the thermal conductivity, density, and specific heat of the metal were taken to be constants, and the resulting problems were linear. However, the conductivity and specific heat are, in fact, functions of temperature.

During the heat transfer process from a single gun firing, the temperature of the metal near the bore surface can rise more than 500 K. The specific heat and thermal conductivity will vary significantly over such a large temperature interval. According to Fourier's heat flux principle, we expect an increase in the conductivity with temperature to augment the resulting heat transfer, and hence further increase the temperature over a scenario in which the conductivity is constant. On the other hand, an increase in specific heat with temperature will have the opposite effect—it will tend to diminish the subsequent temperature rise. Thus, it is difficult to discern, a priori, what effect temperature-dependent thermal properties will have on gun tube heating. We propose to answer this question by solving the nonlinear heat transfer problem with a finite difference solution, using the Crank-Nicolson implicit scheme with iteration. Application of the Kirchhoff transformation, described in Section 3, will render the problem more tractable.

We choose to apply the model to gun barrel heating in tank guns, specifically, the 120-mm M256 cannon for the M1A1 tank. A tank gun is ammunition-limited by the onboard stowage restrictions. This imparts an upper limit to the firing-induced barrel temperature change for tank guns that does not hold for unlimited-fire guns, such as small-caliber chain guns or large-caliber artillery guns. Consequently, we do not treat the broadest possible temperature range for all types of guns; nevertheless, our findings cover a substantial temperature span. The findings indicate that the inclusion of temperature dependence in the modeled thermal properties does not have a significant impact on the barrel temperature predictions over the range investigated.

## 2. THE HEATING MODEL

We shall compute the gun barrel heating for multiple firings by an extension of the model used in Gerber and Bundy (1991). The following assumptions apply here:

(1) Temperature gradients in the axial direction are neglected in comparison to those in the radial direction.

(2) Temperature is axisymmetric in the plane normal to the bore axis. This implies axisymmetric heat input, as well as the neglect of gravity, barrel thickness variation, and other effects that would cause azimuthal dependence.

(3) Feedback of barrel heating to flow in the gun bore is neglected, so that the same bore temperature and convective heat transfer coefficient histories (for a single round) furnish the input data for every round calculated.

(4) Friction heating is neglected.

(5) Thermal expansion of the barrel is not considered to have an effect on the heat transfer process.

(6) The density,  $\rho$ , of the gun barrel metal is constant ( $= 7,827 \text{ kg/m}^3$ ).

(7) The thermal conductivity,  $k$ , and the specific heat,  $c_p$ , of the metal are functions of the temperature,  $T$ , as shown in Figures 1 and 2, respectively. (Figure 1 is based on data from the Material Properties Data Center [1973]; Figure 2 is based on data provided by Joseph Cox, Benet Weapons Laboratory [1990].) Also indicated in the figures are the constant values used by the authors for  $k$  and  $c_p$  in a previous study (Gerber and Bundy 1991):  $k = 38.07 \text{ J/(m s K)}$  and  $c_p = 469.05 \text{ J/kg K}$ . The particular constants chosen are within the range of values spanned in the course of a typical firing scenario.

The abrupt changes in the curve of Figure 2 are associated with the phase transition of steel from a body-centered crystal lattice structure (ferrite) to a face-centered lattice structure (austenite). Even though we adjust the  $k$  and  $c_p$  to account for phase transition in the conduction equation, we shall not (in this report) further modify the conduction equation to account for the latent heat effects of the phase transition.

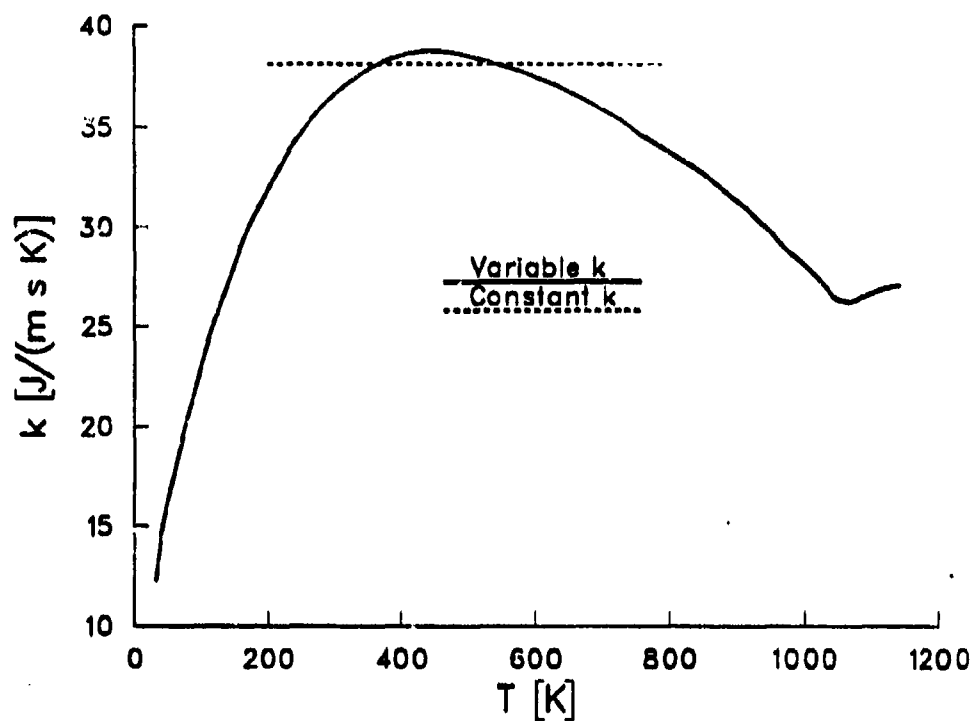


Figure 1. Gun Barrel Thermal Conductivity Variation With Temperature.

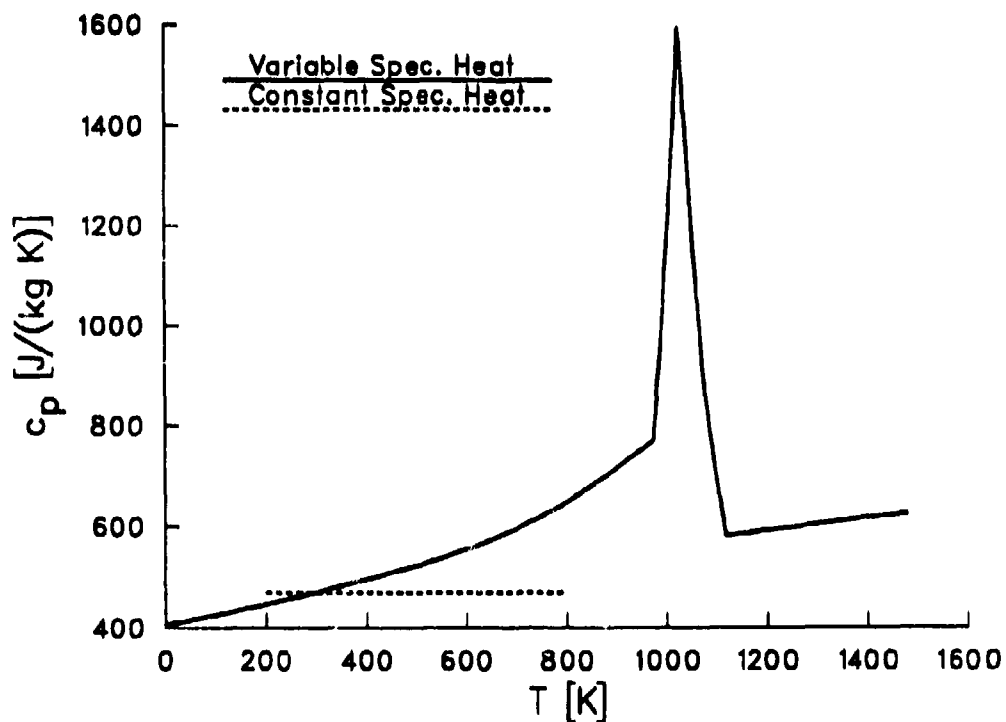


Figure 2. Gun Barrel Specific Heat Variation With Temperature.

### 3. FORMULATION OF THE PROBLEM

**3.1 Statement of the Problem.** We state our problem in terms of cylindrical coordinates:  $r$ ,  $\theta$ , and  $z$ . The radial coordinate,  $r$ , is zero on the axis of the gun tube ( $z$ -axis), and varies from  $R_i$  to  $R_o$ , the radii of the concentric inner and outer walls, respectively (Figure 3). As implied in Section 2, the azimuthal angle,  $\theta$ , does not enter the problem. The axial coordinate,  $z$ , is taken to be zero at the gun's breech. The barrel temperature,  $T(r,z,t)$ , where  $t$  is time measured from the initiation of the first round, is determined by the following differential equation of heat conduction for a stationary, homogeneous, isotropic solid with no internal heat generation (Özisk 1968, 353, or see ADDENDUM):

$$\rho c_p \partial T / \partial t = \text{div} [k \text{ grad } T] .$$

Under the assumptions made in Section 2, this equation reduces to

$$(1/\alpha) \partial T / \partial t = [\partial^2 T / \partial r^2 + (1/r) \partial T / \partial r] + (dk/dT) (\partial T / \partial r)^2 / k , \quad (1)$$

where  $\alpha = k/(\rho c_p)$  is the thermal diffusivity, now a function of  $T$  as shown in Figure 4.

The initial and boundary conditions do not change in form from those for constant properties. Let  $T_\infty$  designate the ambient temperature of the atmosphere, assumed to be constant. The initial condition is

$$T(r) = T_\infty , \quad t = 0, R_i \leq r \leq R_o \quad (z = \text{const}) . \quad (2)$$

The boundary condition at the inner wall is Newton's law of cooling:

$$-k \partial T / \partial r = h_g (T_g - T) , \quad r = R_i , \quad t > 0 \quad (z = \text{const}) , \quad (3)$$

where  $T_g(t,z)$  is the cross-sectional average temperature of the flow in the bore, and  $h_g(t,z)$  is the coefficient of heat transfer between the gas-particle mixture in the bore and the inner wall of the barrel.  $T_g(t,z)$  and  $h_g(t,z)$  are known from interior ballistic computations and thus constitute input.

The boundary condition at the outer wall, which includes both convective and radiative cooling (Özisk 1968; Equations 1-28 and 8-137c), is

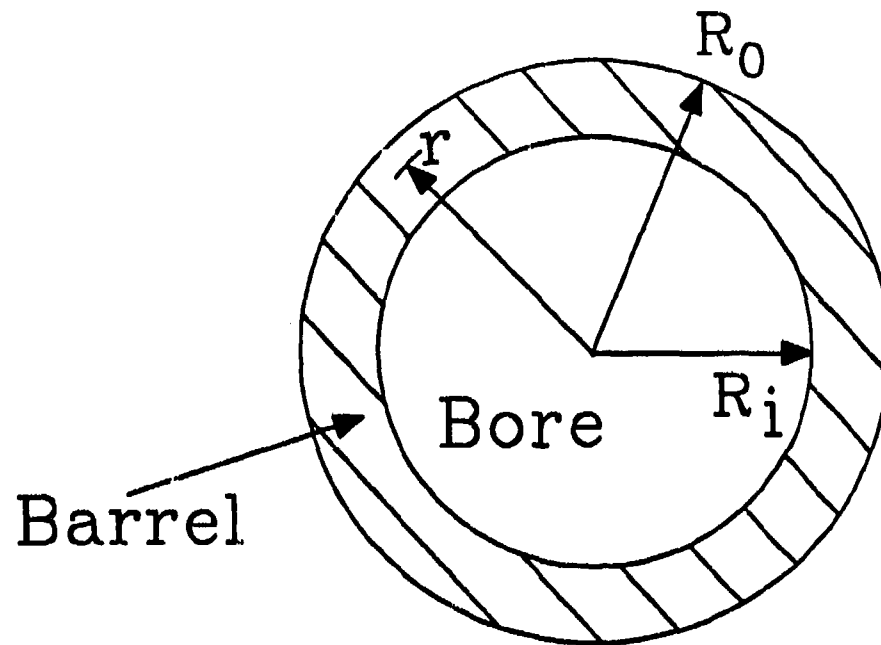


Figure 3. Transverse Cross Section of Gun Barrel.

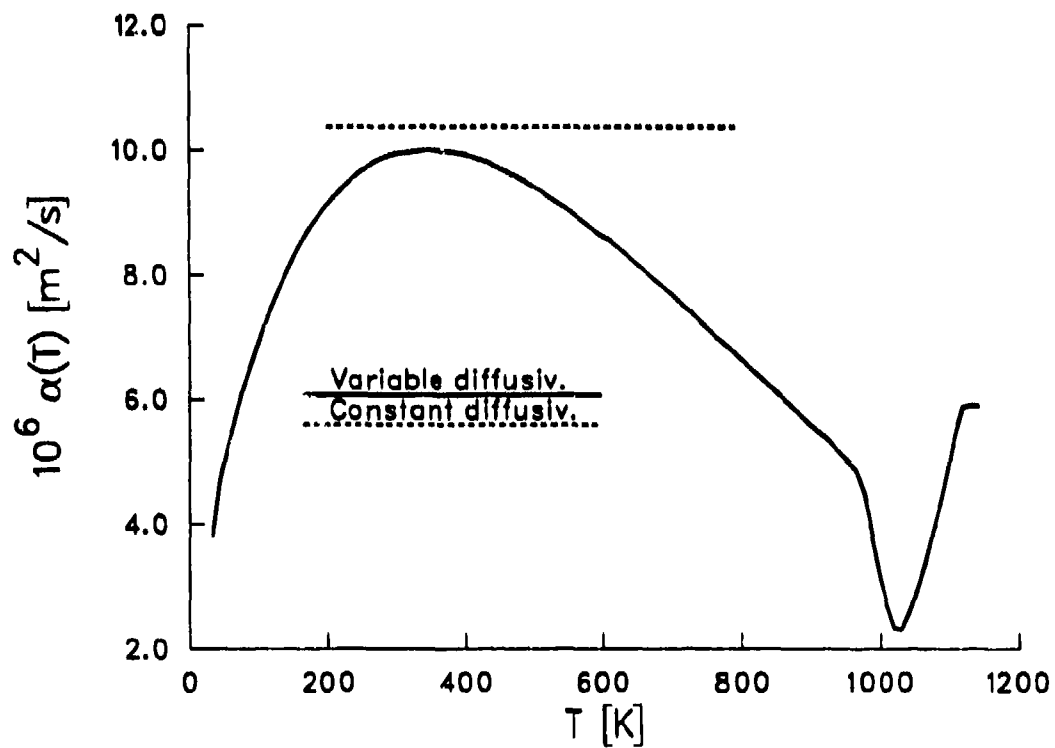


Figure 4. Gun Barrel Thermal Diffusivity Variation With Temperature.

$$-k \partial T / \partial r = h_{\infty} (T - T_{\infty}) + F\sigma (T^4 - T_{\infty}^4) \quad r = R_o, t > 0 (z = \text{const}), \quad (4)$$

where  $h_{\infty} = h_{\infty}(z)$  is the coefficient of convective heat transfer between the barrel wall and the surrounding atmosphere.  $F$  is the radiation interchange factor between the barrel outer wall and the environment (in our case, assumed = 0.95), and  $\sigma$  is the Stefan-Boltzmann constant [ $= 5.669 \times 10^{-8} \text{ J/(m}^2 \text{ s K}^4)$ ].

**3.2 Kirchhoff Transformation.** Equations 1, 3, and 4 show that the heat conduction problem is nonlinear in  $T$ . Introducing the Kirchhoff transformation (Özisk 1968, 353; Boley and Weiner 1960, 141),

$$U(T) = \int_{T_1}^T [k(\tau)/k_1] d\tau, \quad (5)$$

will simplify the conduction equation and boundary conditions.  $U$  has the dimensions of temperature. Here,  $T_1$  is an arbitrarily chosen lower integration limit, namely, a data point in our table for  $k$ ;  $T_1 = 33.4512 \text{ K}$  and  $k_1 = k(T_1) = 12.801 \text{ J/(m s K)}$ . We shall solve the problem posed in terms of  $U$  and then evaluate  $T(U)$  from the inverse of the Kirchhoff transformation.

The variable coefficients  $k(T)$ ,  $c_p(T)$ , and  $\alpha(T)$  are available in tabular form. For a particular argument, the function value is evaluated by 3-point interpolation.  $U(T)$  is obtained in tabular form by numerical integration; it is the solid curve in Figure 5.\* The problem now has the form

$$(1/\alpha) \partial U / \partial t = \partial^2 U / \partial r^2 + (1/r) \partial U / \partial r \quad (6a)$$

$$-k_1 \partial U / \partial r = h_s [T_s - T(U)] \quad r = R_i, t > 0 \quad (6b)$$

$$-k_1 \partial U / \partial r = h_{\infty} (T - T_{\infty}) + F\sigma (T^4 - T_{\infty}^4) \quad r = R_o, t > 0 \quad (6c)$$

$$U = U_{\infty} = U(T_{\infty}) \quad t \leq 0, R_i \leq r \leq R_o \quad (6d)$$

\* It was necessary to extend the tables of  $k$  and  $U$  in Figures 1 and 5 by extrapolation to higher temperatures (2,000 K) because of extremely high surface temperatures attained briefly in the chamber.



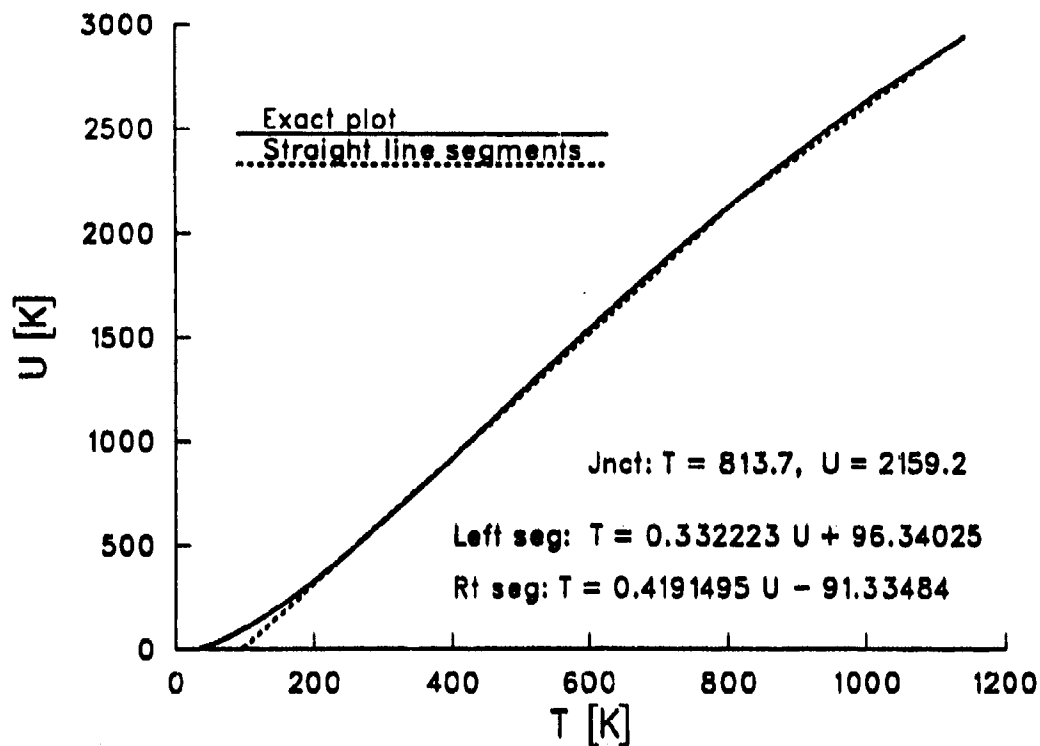


Figure 5. Kirchhoff Transform, U. vs. Temperature.

The diffusivity,  $\alpha = \alpha[T(U)]$ , is a function of  $U$ , shown in Figure 6; thus, Equation 6a is still nonlinear. However, the highly nonlinear term containing  $k$  and its derivative no longer appears. In order to render the finite-difference algorithms for the boundary conditions linear, we introduce two linearizing approximations.

For the first approximation, let  $t^m$  denote the  $m$ th time step of the calculation, when the solution is known, and  $t^{m+1}$  the next time step, when the solution is to be found. At  $r = R_0$  we assume that

$$|T(R_0, t^{m+1}) - T(R_0, t^m)| \ll T(R_0, t^m).$$

Then Equation 6c, at  $r = R_0$  and  $t = t^{m+1}$ , is approximated by

$$k_1 \partial U / \partial r + [h_\infty + 4F\sigma (T^m)^3] T^{m+1} = h_\infty T_\infty + F\sigma [T_\infty^4 + 3(T^m)^4]. \quad (7)$$

This is a reasonable approximation because the temperature changes most slowly at the outer wall; furthermore, the time increment  $\Delta t = t^{m+1} - t^m$  can be decreased to ensure the above inequality condition.

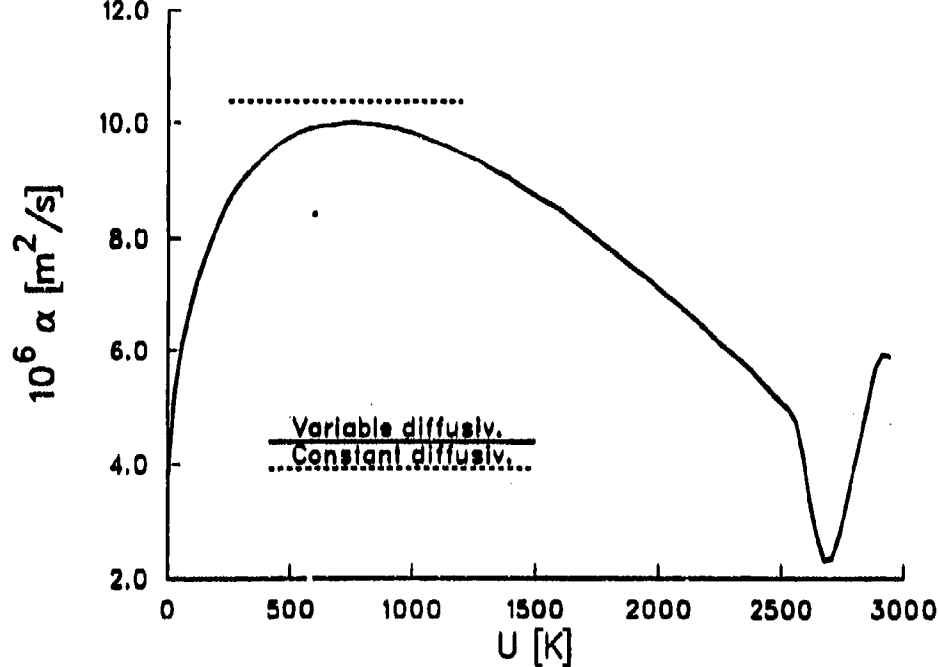


Figure 6. Thermal Diffusivity vs. U.

The second approximation consists of expressing the  $T$  in Equations 6b and 7 as a linear function of  $U$ . The dashed curve in Figure 5 shows that this can be accomplished with reasonable accuracy by the use of two straight line segments. Thus  $T = b_1 U + b_2$ , where

$$b_1 \begin{cases} = 0.332223 & \text{if } U \leq 2,159.2 \\ = 0.4191495 & \text{if } U > 2,159.2 \end{cases} \quad b_2 \begin{cases} = 96.34025 & \text{if } U \leq 2,159.2 \\ = -91.33484 & \text{if } U > 2,159.2 \end{cases}$$

Since it is necessary to choose one of the two line segments a priori at time  $t = t^{m+1}$ , we base the choice on the  $U$  (inner or outer wall) at time  $t = t^m$ .

**3.3 Transformed Radial Coordinate.** We introduce a transformation  $r = r(\xi)$  for  $0 \leq \xi \leq 1$  (Gerber and Bundy 1991) so that the constant increment  $\Delta \xi$  will cluster the nodal points closely together near the inner wall, where  $T$  and  $U$  gradients are the largest. We define the transformation in the following two steps:

$$\begin{cases} \zeta(\xi) = \gamma \xi + (1 - \gamma) \xi^\beta, & (0 \leq \gamma \leq 1, \beta > 2), \\ r = D \zeta + R_i \end{cases} \quad (8)$$

where  $D = R_o - R_i$ , and  $\gamma$  and  $\beta$  are chosen constants. We chose values of  $\gamma = 0.092$  and  $\beta = 2.25$ , which provided a suitable distribution of nodes. Note that  $r = R_i, R_o$  correspond to  $\xi = 0, 1$ , respectively. The actual computations are then carried out in the  $\xi, t$  space; a restatement of the problem in  $\xi$  and  $t$  is provided in Appendix A.

#### 4. INPUT DATA

A detailed discussion of the input to the computations is given in Gerber and Bundy (1991). Briefly, however,  $T_g$  is computed at chosen stations along the bore from the NOVA code (Gough 1980) and  $h_g$  is computed from the Veritay code (Chandra and Fisher 1989a, 1989b), which uses  $T_g$  and other NOVA variables to determine  $h_g$  by the method of Stratford and Beavers (1961).

Figures 7a and 7b show representative  $T_g$  and  $h_g$  histories at two stations on a 120-mm M256 gun barrel. It is seen that  $T_g$  and  $h_g$  remain constant until the base of the projectile passes the given station at time  $t = t_g$ . At this time, these variables rise suddenly to a peak, then they decrease gradually, with  $h_g$  decaying significantly faster than  $T_g$ .

All the computations reported here were performed for a 120-mm M256 tank gun firing a DM13 round.\* The density of the gun barrel metal is taken to be  $\rho = 7,827.0 \text{ kg/m}^3$ . The thermal conductivity and specific heat are supplied by tables that are plotted in Figures 1 and 2, respectively. The value for  $h_{\infty} = 6.0 \text{ J/(m}^2 \text{ s K)}$  was obtained from experiments conducted by Bundy on a shrouded M256 barrel.

#### 5. FINITE-DIFFERENCE CALCULATION

For the finite-difference calculations, the interval  $0 \leq \xi \leq 1$  (corresponding to  $R_i \leq r \leq R_o$ ) is divided by equally spaced nodal (or grid) points into  $NI$  subintervals. The constant  $\xi$  increment is  $\Delta\xi = 1/NI$ , and the location of the nodes is given by  $\xi_j = (j - 1) \Delta\xi$  ( $j = 1, 2, \dots, NI + 1$ ). Derivatives at node  $j$  are approximated as follows:

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\* Table 1 of Gerber and Bundy (1991) describes the shape of the gun barrel.

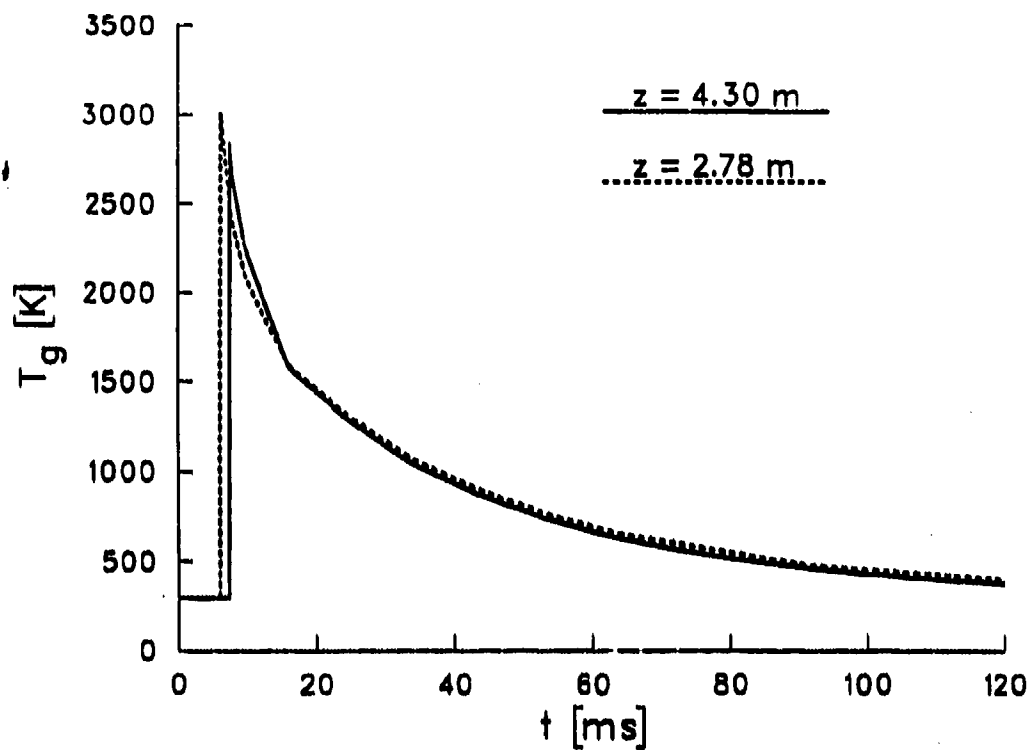


Figure 7a. Bore Gas Temperature Histories at Two Axial Locations.

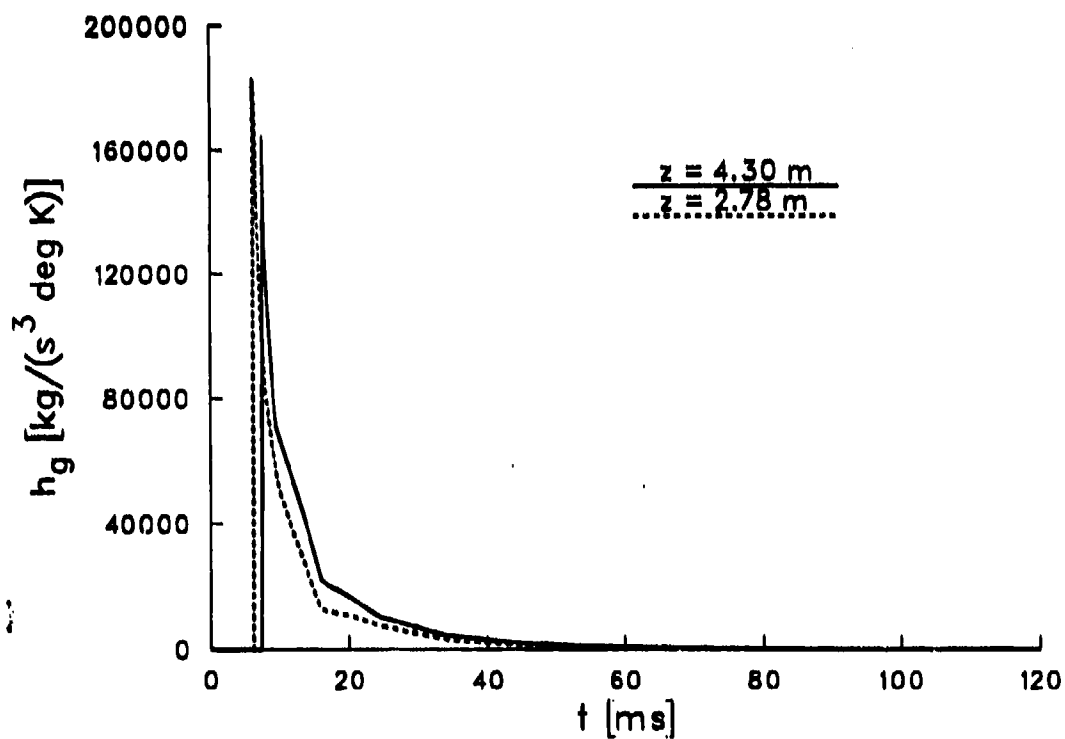


Figure 7b. Convective Heat Transfer Coefficient Histories at Inner Wall of Gun Barrel at Two Axial Locations.

$$(\partial U / \partial \xi_j) = (-3 U_j + 4 U_{j+1} - U_{j+2}) / (2 \Delta \xi) \quad (j = 1) \quad (9a)$$

$$(\partial U / \partial \xi_j) = (U_{j+1} - U_{j-1}) / (2 \Delta \xi) \quad (j = 2, \dots, NI) \quad (9b)$$

$$(\partial U / \partial \xi_j) = (U_{j-2} - 4 U_{j-1} + 3 U_j) / (2 \Delta \xi) \quad (j = NI + 1) \quad (9c)$$

$$(\partial^2 U / \partial \xi_j^2) = (U_{j-1} - 2 U_j + U_{j+1}) / (\Delta \xi)^2 \quad (j = 2, \dots, NI). \quad (9d)$$

If we let the time increment be  $\Delta t = t^{n+1} - t^n$ , then, in the Crank-Nicolson scheme employed here (Özisk 1968, 402) to obtain the solution at time  $t = t^{n+1}$  ( $j = 2, \dots, NI$ ),

$$(U_j^{n+1} - U_j^n) = (\Delta t / 2) [(\partial U / \partial t)^n + (\partial U / \partial t)^{n+1}]_j. \quad (10)$$

By Equation A-4a, Equation 10 becomes

$$U_j^{n+1} - (\Delta t / 2) H_j^{n+1} = U_j^n + (\Delta t / 2) H_j^n, \quad (11)$$

where

$$H = (\alpha / D^2) [f_1(\xi) \partial^2 U / \partial \xi^2 + f_2(\xi) \partial U / \partial \xi],$$

and  $f_1(\xi)$  and  $f_2(\xi)$  are defined in Equations A-3.

The finite-difference approximations to the equations and boundary conditions are produced by substituting the derivative approximations of Equation 9 into Equations 11, A-4b, and A-4c, and then collecting the terms. The following set of equations for  $U_j^{n+1}$  may then be written:

$$\sum_{j=1}^{NI+1} A_{nj} U_j^{n+1} = d_n \quad (n = 1, 2, \dots, NI+1). \quad (12)$$

The coefficients  $A_{nj}$  and  $d_n$  are given in Appendix B. The  $d_n$ 's involve  $U$  values for the previous timestep  $t = t^n$ . In most cases, we have used  $NI = 100$ .

The coefficients  $A_{nj}$  ( $j = 2, \dots, NI$ ) are linear functions of  $\alpha(U)$  in Equations (B-3) in Appendix B, and are not predetermined constants. A successive approximation procedure is used to solve Equation 12. First, an initial estimate is made for the  $\alpha$ 's, e.g.,  $(\alpha_j^{m+1})^{1st} = \alpha_j^m$ , which are known. These  $\alpha$ 's, when substituted into the  $A_{nj}$ 's, make Equation 12 a linear system, which is solved for  $(U_j^{m+1})^{1st}$  by a standard FORTRAN routine. The next  $\alpha$  approximation is  $(\alpha_j^{m+1})^{2nd} = \alpha[(U_j^{m+1})^{1st}]$ . Then  $(U_j^{m+1})^{2nd}$  is the solution to Equation 12 when the  $(\alpha_j^{m+1})^{2nd}$  are substituted into the  $A_{nj}$ 's; and  $(\alpha_j^{m+1})^{3rd} = \alpha[(U_j^{m+1})^{2nd}]$ . This procedure is repeated until convergence is obtained; i.e.,

$$|(U_j^{m+1})^{(i+1)th} - (U_j^{m+1})^{ith}| < C_{conv} \quad (j = 1, 2, \dots, NI + 1),$$

where  $C_{conv}$  is a chosen arbitrarily small positive constant. Three iterations have been required, on the average, for our calculations.

The computer program contains a subroutine prescribing  $\Delta t$  as a function of  $t$  within a firing cycle (see Appendix C). The  $\Delta t$  is made sufficiently small early in the cycle to resolve the highly transient phenomena; then it is increased to minimize computation time. The Crank-Nicolson method is stable for all values of  $\Delta t$ , and there are no restrictions on the relative sizes of  $\Delta t$  and  $\Delta x$ .

After the problem is solved for  $U$ ,  $T$  is found by inverting the Kirchhoff transformation. This step is accomplished by applying three-point interpolation in the stored  $T$ ,  $U$  table.

## 6. ACCURACY CHECK

The condition of energy balance can be employed as an accuracy check. We define two integral functions: (1)  $Q_F$  = net energy per unit length that has entered and exited the barrel since  $t = 0$ , and (2)  $Q_b$  = change in barrel heat content per unit length since  $t = 0$ .

$$\begin{aligned} Q_F = & 2\pi R_i \int_0^1 h_s(\tau) [T_s(\tau) - T(R_i, \tau)] d\tau - 2\pi R_o \int_0^1 h_w [T(R_o, \tau) - T_w] d\tau \\ & - 3.5619 \times 10^{-7} FR_o \int_0^1 [(T(R_o, \tau))^4 - T_w^4] d\tau \quad J/m \end{aligned} \quad (13)$$

$$Q_b = 2 \pi \rho \int_{R_i}^{R_o} \left[ \int_{T_m}^{T(r)} c_p(\bar{T}) d\bar{T} \right] r dr \quad J/m \quad (14)$$

A necessary, but not sufficient, condition for accuracy is that  $Q_b = Q_r$  at each time step.

## 7. COMPUTATIONS

Computations were performed to determine the heating of a 120-mm M256 gun tube firing DM13 rounds. We chose just two of the many possible firing scenarios. The first chosen firing sequence consisted of 16 rounds at 40-s intervals, followed by 5 rounds at 30-s intervals. Figure 8a compares the two outer wall temperature histories at  $z = 4.30$  m, about 1 m from the muzzle. At the end of the firing sequence ( $t = 780$  s) there is a 3.8 K discrepancy between the two computations. This represents about 2.6% of the total outer wall barrel temperature change (total temperature change is defined as  $T[t] - T_m$ ). A full-scale plot in  $T$  vs.  $t$  at the inner wall would not reveal the differences between linear and nonlinear output. Figure 8b shows a greatly magnified picture of the histories at one of the inner wall temperature peaks, where the discrepancy maxima occur. The difference at this point is about 25 K, or about 3.4% of the inner wall barrel temperature change.

The main comparisons were made for our second firing sequence, namely, the one shown in Table 1, which is almost identical to Scenario #2 in Table 1 of Artus and Hasenbein (1989). (Our final burst is seven rounds; that of Scenario #2 is six rounds.) This scenario represents the case where all the M1A1 tank-stowed ammunition is fired as fast as possible, with two intermediate cooldown periods corresponding to accessing different ammunition storage compartments. We have computed the change for 41 rounds, whereas the M1A1 tank stows 40 rounds. Also, we have assumed that all 41 rounds are DM13 (kinetic energy [KE] rounds), whereas in actuality the M1A1 stores a mix of KE and HEAT (high-explosive antitank) rounds.

The simulations were performed at two locations on the barrel, 4.30 m and 0.615 m from the breech, for a high and a low ambient temperature: 322 K (120° F) and 233 K (-40° F), respectively. These two locations correspond to positions where the barrel is relatively thin and thick, respectively. The two ambient temperatures correspond to frequently used Army hot and cold testing conditions. The initial gun barrel temperature was the same as that of the surroundings.

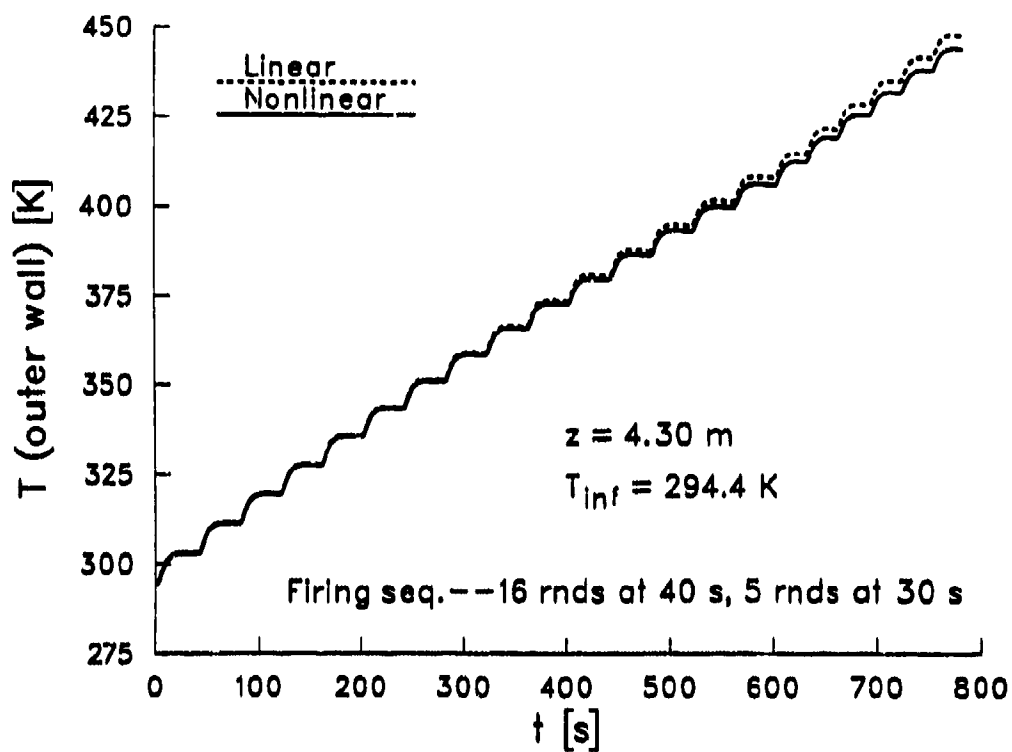


Figure 8a. Linear and Nonlinear Outer Wall Temperature Histories.

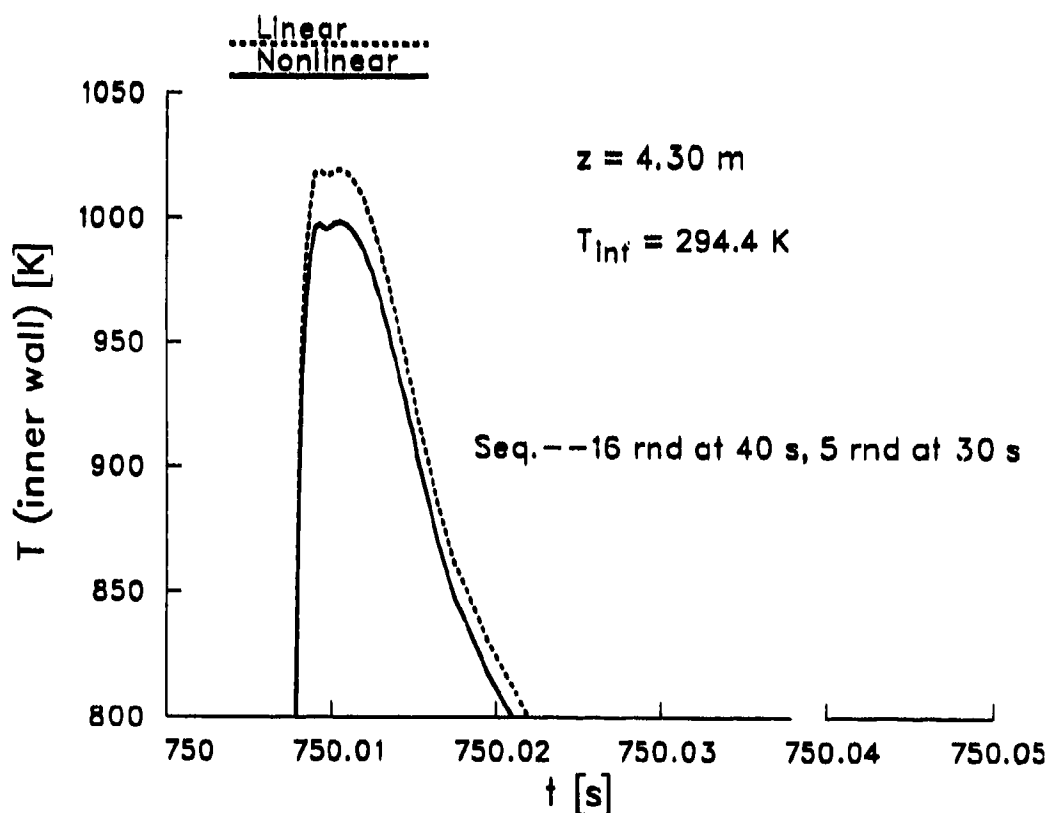


Figure 8b. Detailed Comparison of Linear and Nonlinear Calculations at an Outer Wall Temperature History Peak.



Table 1. Firing Sequence

- |                              |
|------------------------------|
| 1. 17 rounds at 7 rounds/min |
| 2. 5 min cool down           |
| 3. 17 rounds at 7 rounds/min |
| 4. 5 min cool down           |
| 5. 7 rounds at 7 rounds/min  |
| 6. Cool down                 |

Figure 9a presents a comparison of the two inner wall heating predictions at  $z = 4.30$  m for  $T_{\infty} = 322$  K. (For convenience,  $T(R_i)$  has been truncated at 600 K.) The largest difference occurs at the end of each firing burst; at  $t = 1,000$  s, for example, the evaluations differ by about 4.8% of the total temperature change, similar to the results from our first firing sequence. Figure 9b shows a greatly magnified picture at one of the temperature peaks. The relative difference between maximum values is about 3.3% of the total temperature change.

In Figure 10, the outer wall temperature calculations are compared for the two ambient temperatures at  $z = 4.30$  m. At  $t = 1,000$  s, shortly after the end of the final burst, differences are 4.5% and 2.5% of total temperature change for  $T_{\infty} = 322$  K and  $T_{\infty} = 233$  K, respectively.

The next two figures show heating results at  $z = 0.615$  m, close to the breech, for the scenario of Table 1. In Figure 11, the linear and nonlinear calculations are seen to be almost identical, more so even than at  $z = 4.30$  m. However, the inclusion of temperature dependence results in a slightly higher, rather than lower (as in Figure 10), outer wall temperature prediction. Figure 12 depicts the inner wall temperature histories under cold and hot firing conditions; the linear and nonlinear curves are almost coincident, at least for the cooling periods.

Figure 13 shows comparisons for the spatial variation of temperature from inner to outer wall for  $T_{\infty} = 233$  K. At  $z = 0.615$  m, the barrel thickness is 95 mm, compared to 23 mm at  $z = 4.30$  m, and the temperature is not able to equilibrate as thoroughly during the 5-minute cooldown periods as it does at the thinner cross section. This behavior is reflected in the fact that  $T$  at the thicker section is still varying significantly from the inner to the outer wall. Also, note that the differences between the linear and nonlinear predictions are noticeably greater for thinner barrel cross sections than for the thicker cross sections.

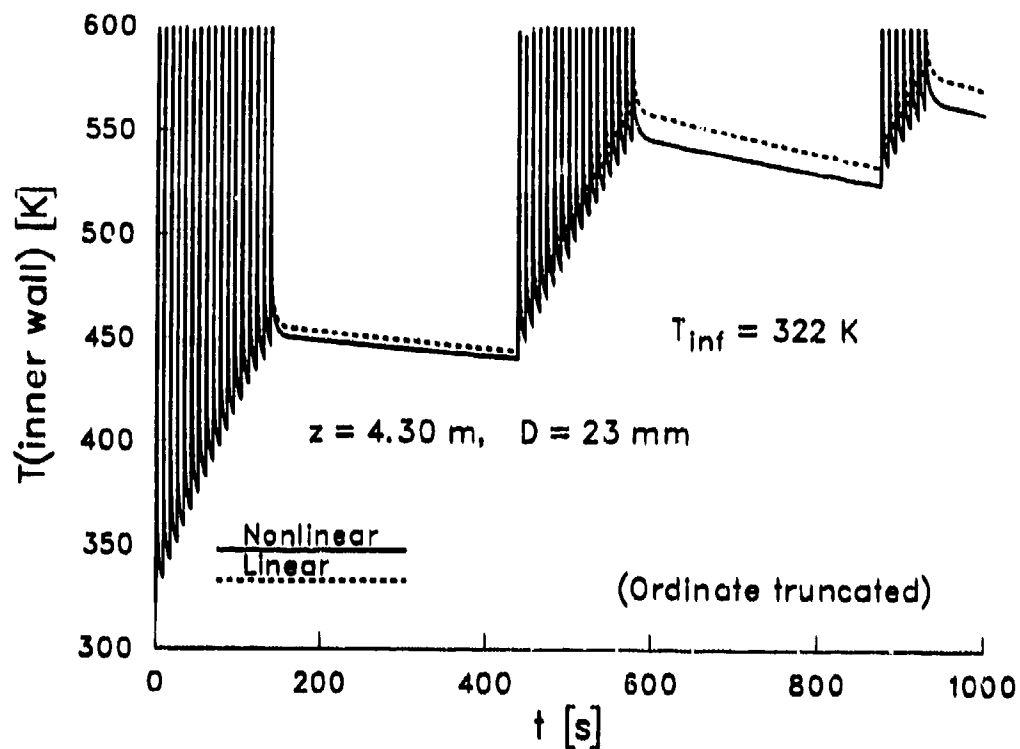


Figure 9a. Comparison of Linear and Nonlinear Inner Wall Temperature Histories at  $z = 4.30$  m (Firing Sequence of Table 1).

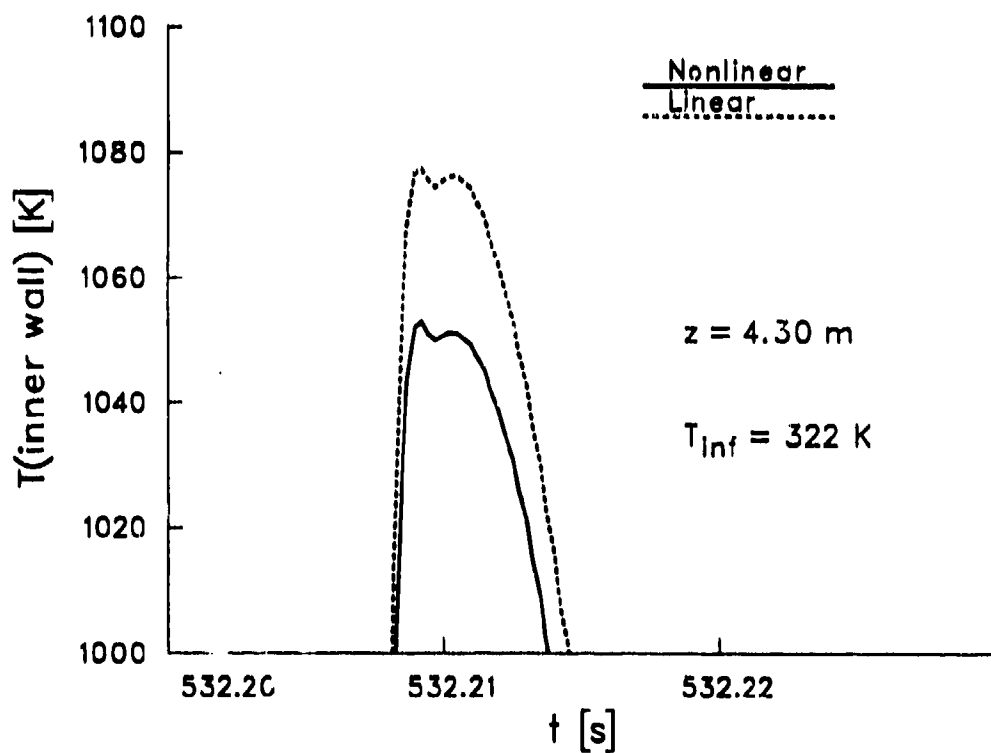


Figure 9b. Comparison of Linear and Nonlinear Calculations at an Inner Wall Temperature History Peak at  $z = 4.30$  m (Firing Sequence of Table 1).

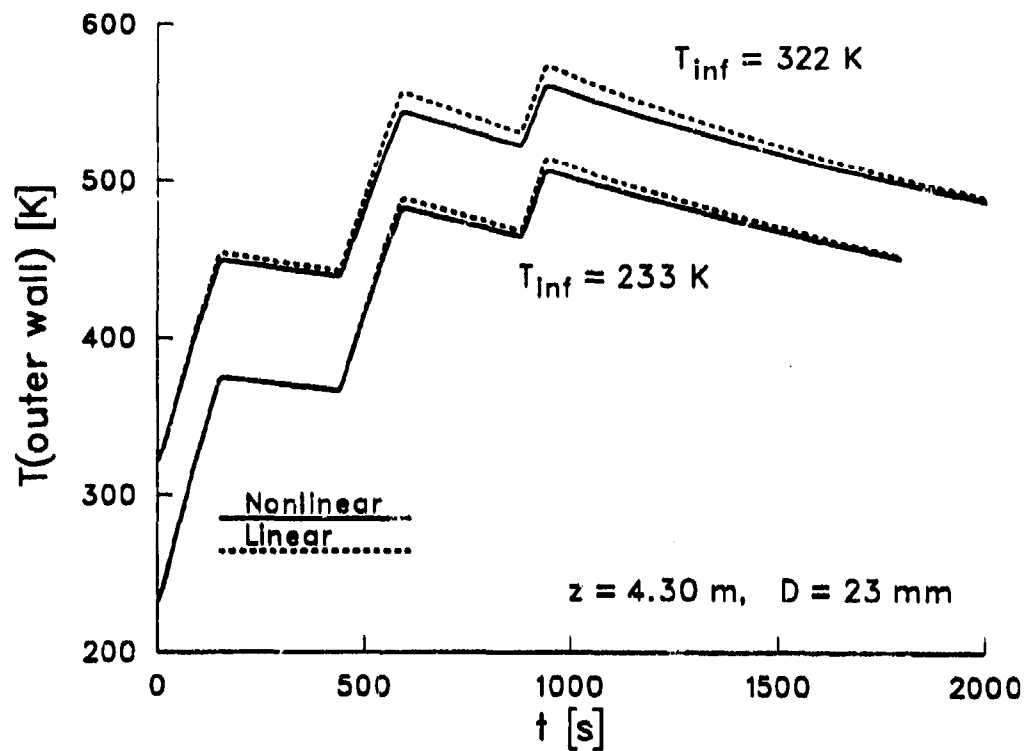


Figure 10. Comparison of Linear and Nonlinear Outer Wall Temperature Histories at  $z = 4.30$  m (Firing Sequence of Table 1).

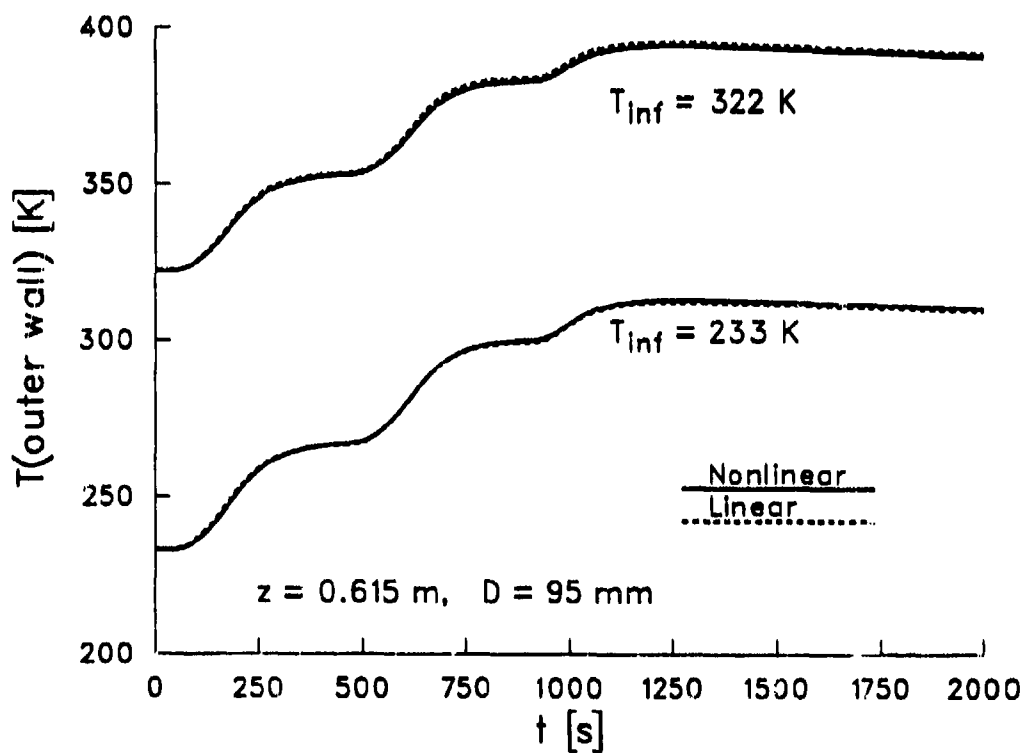


Figure 11. Comparison of Linear and Nonlinear Outer Wall Temperature Histories at  $z = 0.615$  m (Firing Sequence of Table 1).

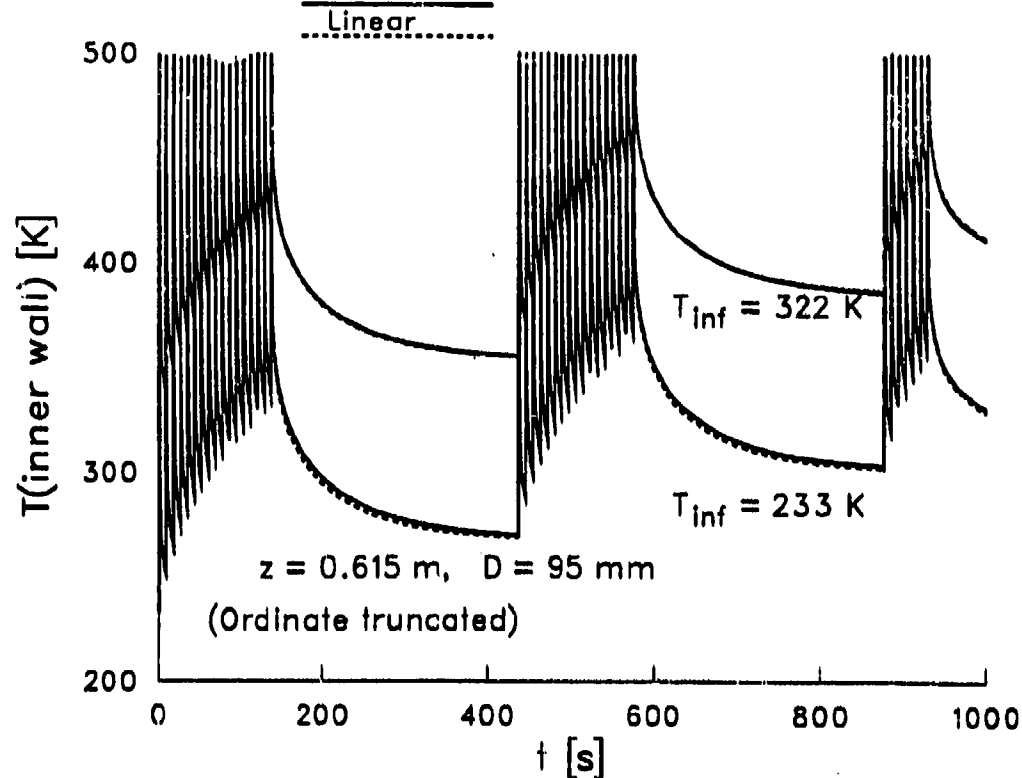


Figure 12. Comparison of Linear and Nonlinear Inner Wall Temperature Histories at  $z = 0.615 \text{ m}$  (Firing Sequence of Table 1).

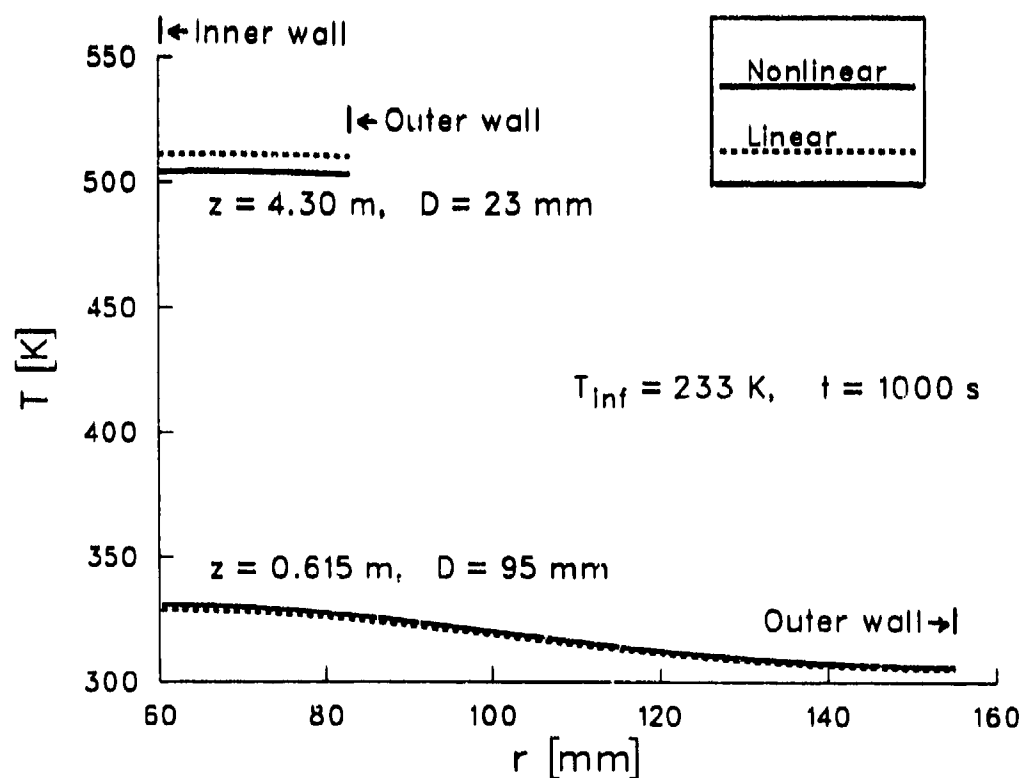


Figure 13. Comparison of Linear and Nonlinear Radial Temperature Profiles at  $t = 1.000 \text{ s}$  (Firing Sequence of Table 1).

## 8. DISCUSSION

This study investigated the "error" introduced into gun barrel heating calculations by assuming constant values of thermal conductivity and specific heat for the metal (linear model) vs. incorporating temperature dependence into these thermal properties (nonlinear model). The constants  $k$  and  $c_p$  chosen are within the range spanned in the course of a typical firing scenario. Figures 1 and 2 show that the constant values differ significantly from the more accurate temperature-dependent  $k$  and  $c_p$  at very high temperatures, say above 800 K. For tank guns, however, limitations on the firing rate and number of rounds that can be fired greatly restrict the duration of extremely high temperatures, so that drastic differences need not occur between predictions of the linear and nonlinear models.

It is not feasible to conduct parameter studies for a wide variety of parameters and firing scenarios. We have concentrated on what we believe to be the two cases that should show the greatest difference between these two models: the rapid expenditure of all tank-stowed ammunition under very hot and very cold initial barrel temperatures. We draw several inferences from our limited computer runs:

- (1) Use of variable thermal properties can either increase or decrease the barrel temperature (e.g., Figure 13), depending on the case (that is, on the totality of the computer input parameters, including  $T_b$  and  $h_b$  histories).
- (2) Differences between linear and nonlinear model temperatures are greater for thinner barrel cross sections (e.g., Figures 9a, 12, and 13). This may be due to the more fundamental fact that thinner cross sections are cycled through a greater temperature range; as a result, the effect of temperature-dependent thermal properties is larger.
- (3) The differences between the outputs of the two models, constant and variable thermal properties, are below 5% for the cases considered. For most tank gun applications, the constant value (or linear) model should be acceptable. In fact, it might be preferred, since it requires roughly one-third of the CPU time of the nonlinear model in each computer run.

## 9. ADDENDUM: HEAT CONDUCTION EQUATION

In one dimension the heat conduction equation can be understood as follows. For a solid material body, the heat entering per unit area, per unit time through the y-z plane at x, for example, will be given by:

$$Q_{in} = -k \frac{\partial T}{\partial x} \Big|_x .$$

Likewise, the heat leaving the through the y-z plane at  $x+dx$  will be:

$$Q_{out} = -k \frac{\partial T}{\partial x} \Big|_{x+dx} .$$

The change in internal energy per unit time, per unit volume, will be:

$$\frac{du}{dt} = \frac{Q_{in} - Q_{out}}{dx} = \frac{-k \frac{\partial T}{\partial x} \Big|_x + k \frac{\partial T}{\partial x} \Big|_{x+dx}}{dx} .$$

Expanding the temperature gradient at  $x + dx$  about x, retaining terms up to first order in dx, then substituting yields:

$$\frac{du}{dt} = \frac{\partial (k \frac{\partial T}{\partial x})}{\partial x} .$$

From thermodynamics, if the elements boundaries are fixed, then we are assured that  $u = u(T)$ , and

$$\frac{du(T)}{dt} = \frac{du(T)}{dT} \frac{\partial T}{\partial t} = \rho c_p \frac{\partial T}{\partial t} .$$

Hence, equating the last two expressions, yields:

$$\rho c_p \frac{\partial T}{\partial t} = \frac{\partial (k \frac{\partial T}{\partial x})}{\partial x} .$$

If these arguments were repeated in three dimensions, the results would lead to the heat conduction equation of Section 3.1, viz.:

$$\rho c_p \frac{\partial T}{\partial t} = \nabla \cdot (k \nabla T) .$$

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**APPENDIX A:**  
**STATEMENT OF PROBLEM IN  $\xi, t$  VARIABLES**

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We repeat the transformation of Equation 8:

$$\zeta(\xi) = \gamma\xi + (1 - \gamma)\xi^\beta \quad (0 < \gamma \leq 1, \beta > 2) \quad r = D\zeta + R_1. \quad (\text{A-1})$$

Then

$$d\zeta/d\xi = \zeta' = \gamma + \beta(1 - \gamma)\xi^{\beta-1}, \quad d^2\zeta/d\xi^2 = \zeta'' = \beta(\beta - 1)(1 - \gamma)\xi^{\beta-2}$$

$$\zeta'(0) = \gamma, \quad \lambda_1 = \zeta'(1) = \gamma + \beta(1 - \gamma)$$

$$\lambda_2 = \zeta''(1) = \beta(\beta - 1)(1 - \gamma). \quad (\text{A-2})$$

We define  $f_1$  and  $f_2$ :

$$f_1 = 1/(\zeta')^3, \quad (\text{A-3})$$

$$f_2 = (D/\zeta')/(D\zeta + R_1) - \zeta''/(\zeta')^3.$$

The transformed problem for U is

$$\partial U/\partial t = (\alpha/D^2)[f_1(\xi)\partial^2 U/\partial \xi^2 + f_2(\xi)\partial U/\partial \xi] = H(\xi, t) \quad (\text{A-4a})$$

$$\partial U/\partial \xi = [h_1(t)D\gamma/k_1][b_1 U + b_2 - T_g(t)] \quad \text{at } \xi = 0 \quad (\text{A-4b})$$

$$[k_1/(D\lambda_1)]\partial U/\partial \xi + b_1 w_1 U = w_2 - b_2 w_1 \quad \text{at } \xi = 1. \quad (\text{A-4c})$$

where

$$\left. \begin{aligned} w_1 &= h_- + 4F\sigma(T^-)^3 \\ w_2 &= h_-T_- + F\sigma[T_-^4 + 3(T^-)^4] \end{aligned} \right\} \text{ at } \xi = 1. \quad (\text{A-5})$$

The initial condition is

$$U = U_- = U(T_-) \quad t = 0, 0 \leq \xi \leq 1. \quad (\text{A-6})$$

**APPENDIX B:**  
**COEFFICIENTS OF EQUATION 12**

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We first define  $v_1(t)$  and  $v_2(t)$ :

$$v_1 = 2 \Delta \xi h_s(t) D \gamma / k_1, \quad v_2 = v_1 [b_2 - T_s(t)],$$

where  $b_1$  and  $b_2$  are given in the last paragraph of Section 3.2. Then,

$$\left. \begin{aligned} A_{11} &= -(3 + b_1 v_1), & A_{12} &= 4, & A_{13} &= -1 \\ d_1 &= v_2 \end{aligned} \right\}. \quad (B-1)$$

At the outer wall, where  $j = NP = NI + 1$ , the coefficients of the boundary condition are determined by the following sequence of formulas:

$$w_1 = h_- + 4 F \sigma (T_{NP}^*)^3$$

$$w_2 = h_- T_- + F \sigma [T_-^4 + 3 (T_{NP}^*)^4]$$

$$\lambda_1 = \gamma + \beta(1 - \gamma)$$

$$q_1 = 2 \Delta \xi w_1 b_1 D \lambda_1 / k_1$$

$$q_2 = 2 \Delta \xi D \lambda_1 (w_2 - w_1 b_2) / k_1$$

$$\left. \begin{aligned} A_{NP,NP-2} &= 1, & A_{NP,NP-1} &= -4, & A_{NP,NP} &= (3 + q_1) \\ d_{NP} &= q_2 \end{aligned} \right\}. \quad (B-2)$$

We next define four functions,  $p_{1j}$ ,  $p_{2j}$ ,  $p_{3j}$ , and  $G_j^*$  for  $2 \leq j \leq NI$ :

$$p_{1j} = [f_{1j} / (\Delta \xi)^2] - [f_{2j} / (2 \Delta \xi)]$$

$$p_{2j} = -2 f_{1j} / (\Delta \xi)^2$$

$$p_{3j} = [f_{1j}/(\Delta\xi)^2] + [f_{2j}/(2 \Delta\xi)]$$

$$G_j^m = (\alpha_j^m/D^2) [p_{1j} U_{j-1}^m + p_{2j} U_j^m + p_{3j} U_{j+1}^m],$$

where  $f_{1j} = f_1(\xi_j)$  and  $f_{2j} = f_2(\xi_j)$  are defined in Appendix A. Then

$$A_{j, j-1} = -\Delta t \alpha_j^{m+1} p_{1j}/(2 D^2) \quad (B-3a)$$

$$A_{j, j} = 1 - \Delta t \alpha_j^{m+1} p_{2j}/(2 D^2) \quad (B-3b)$$

$$A_{j, j+1} = -\Delta t \alpha_j^{m+1} p_{3j}/(2 D^2) \quad (B-3c)$$

$$d_j = U_j^m + (\Delta t/2)G_j^m. \quad (B-3d)$$

All other  $A_{ij}$  are equal to zero.



**APPENDIX C:**  
**TIMESCALE**

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Here, time =  $t'$  will refer to time within one firing cycle —  $t' = 0$  at the beginning of the cycle. Six constants are given:  $t_d$ ,  $t_r$ ,  $t'_1$ ,  $t'_2$ ,  $\Delta t'_1$  and  $\Delta t'_2$ . Here,  $t_d$  is the delay time for the rapid rise in  $T_g$  and  $h_g$  from initial conditions, and  $t_r$  is the time between successive firings. The time increment  $\Delta t$  ( $t'$ ) is given by the following function:

$$\Delta t = t_d - \Delta t'_1 \quad 0 \leq t' < t_d - \Delta t'_1$$

$$\Delta t = \Delta t'_1 \quad t_d - \Delta t'_1 \leq t' < t'_1$$

$$\Delta t = C_1 + C_2 t' \quad t'_1 \leq t' < t'_2$$

$$\Delta t = \Delta t'_2 \quad t'_2 \leq t'$$

where

$$C_2 = (\Delta t'_2 - \Delta t'_1) / (t'_2 - t'_1) \text{ and } C_1 = \Delta t'_2 - C_2 t'_2$$

$$(\text{If } t' + \Delta t'_2 > t_r, \text{ set } \Delta t = t_r - t').$$

A typical set of values of the parameters would be the following:

$$t'_1 = 0.018 \text{ s}, \quad t'_2 = 10.0 \text{ s}, \quad \Delta t'_1 = 0.00025 \text{ s}, \quad \Delta t'_2 = 6.0 \text{ s}.$$

These values were chosen on the basis of experience from many empirical studies made with the numerical parameters.

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## LIST OF SYMBOLS

$A_{nj}$	coefficient in Equation 12 for barrel temperature
$b_1, b_2$	constants in $T \approx b_1 U + b_2$ , see last paragraph of Section 3.2
$c_p(T)$	specific heat of gun barrel [J/(kg K)]
$D$	$= (R_o - R_i)$ = thickness of gun barrel [m, mm]
$d_n$	coefficient in barrel temperature equations, right-hand side of Equation 12
$F$	radiation interchange factor between barrel outer wall and environment, Equation 4
$f_1, f_2$	given functions of $\xi$ , Equation A-3
$H$	function defined in Equation A-4
$h_b$	heat transfer coefficient - bore to gun barrel [J/(m <sup>2</sup> s K)]
$h_{\infty}$	heat transfer coefficient - gun barrel to ambient air [J/(m <sup>2</sup> s K)]
$j$	index indicating radial location of a nodal point ( $j = 1$ for inner wall, $j = NP$ for outer wall)
$k(T)$	thermal conductivity of gun barrel [J/(m s K)]
$k_i$	$k(T_i)$
$NI$	number of intervals in $R_i \leq r \leq R_o$ and $0 \leq \xi \leq 1$
$NP$	$= NI + 1$
$Q_b$	increase in heat content per unit length of barrel since $t = 0$ [J/m]
$Q_F$	net quantity of heat per unit length that has entered barrel since $t = 0$ [J/m]
$R_i, R_o$	radial coordinates of inner and outer walls, respectively, of gun barrel [m, mm]
$r$	radial coordinate in transverse plane [m, mm] ( $r = 0$ at axis of gun bore)
$T$	temperature in gun barrel [K]
$T_b, T_{\infty}$	temperatures in the bore and ambient air, respectively [K]
$T_{int}$	$= T_{\infty}$
$T_1$	prescribed lower integration limit in definition of $U$ (Equation 5)

$t$	time from initiation of first round [s, ms]
$t_d$	delay time at given $z$ for rapid rise in $T_g$ and $h_g$ [s, ms]
$t_r$	time interval between two successive rounds [s, ms]
$t'$	time measured within a firing cycle [s, ms]
$U(T)$	Kirchhoff transformation function, defined in Equation 5 [K]
$U_j$	$= U(\xi = [j - 1] \Delta\xi)$
$U_\infty$	$= U(T_\infty)$
$z$	axial coordinate ( $z = 0$ at breech) [m]
$\alpha(T)$	$= k/(\rho c_p)$ , thermal diffusivity of gun barrel [ $m^2/s$ ]
$\beta, \gamma$	prescribed constants in coordinate transformation, Equation 8
$\Delta t$	computational time increment $= (t^{m+1} - t^m)$ [s]
$\Delta\xi$	$= 1/NI$ , constant step size in $\xi$ ( $0 \leq \xi \leq 1$ )
$\zeta$	$= \gamma \xi + (1 - \gamma) \xi^2$
$\theta$	azimuthal coordinate in transverse plane
$\lambda_1, \lambda_2$	$\zeta'(\xi = 1), \zeta'(1)$ , defined in Appendix A
$\xi$	transformed radial variable, Equation 8; independent variable in Equation A-4a
$\rho$	density of gun barrel metal [ $kg/m^3$ ]
$\sigma$	Stefan-Boltzman constant $= 5.669 \times 10^{-8} J/(m^2 s K^4)$

#### Superscript

$m$	index of current time step, when solution is known
$m + 1$	index of next time step, when solution is to be calculated

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